

METHODS OF APPLIED MATHEMATICS COMPREHENSIVE
EXAMINATION AUGUST 2017

Work on as many of the following problems as possible. Turn in *all* your work.

- (1) Consider the projectile problem of a body (mass m), dropped vertically from a point above the Earth's equator (assume a constant gravitational field and that the body falls on a line connecting its drop point to the Earth center). Assuming that the object always experiences a cubic drag force as a function of its velocity:
- Write down the Newtonian dynamics.
 - Non-dimensionalize with respect to viscous scales (terminal velocity and initial height), identify non-dimensional parameters.
 - Evaluate a perturbation solution in the small non-dimensional parameter and sketch the solution
 - For an observer on the Earth's equator, write the leading order trajectory from the observer's perspective (assuming the Earth is in constant rotation with respect to the north-south pole axis).

- (2) Consider the eigenvalue problem on the real line $x \in \mathbb{R}$

$$\epsilon y'' - (U(x) + \lambda)y = 0, \quad y(x) \rightarrow 0, \text{ as } |x| \rightarrow \infty.$$

with the potential:

$$U(x) = a \delta(x - L) + b \delta(x + L)$$

where $\delta(x)$ is the Dirac delta function. First, set $b = 0$ and solve for the eigenvalue(s) and eigenfunction(s), as the parameter a varies over the real line. Second, calculate the eigenvalue(s) and eigenfunctions with when both a and b vary over the real line. How does the single delta function result compare with the double delta case as $L \rightarrow \infty$?

- (3) Consider the boundary value problem

$$\epsilon \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = \epsilon u^2,$$

$$u(0, y) = g(y), \quad u(1, y) = 0, \quad u(x, 0) = f(x), \quad u(x, 1) = 0,$$

in $(x, y) \in [0, 1] \times [0, 1]$, for some functions $f(x)$, $g(y)$. Find the first term in an asymptotic expansion of the solution as $\epsilon \rightarrow 0$. The functions f and g can be taken to be smooth but are otherwise generic (i.e., do not take special values in $[0, 1]$). If you cannot solve any of the differential equations necessary for your procedure, just set them up properly (i.e., state the proper equation with proper initial and/or boundary conditions etc.).

- (4) Consider the functions

$$f(z) = \sqrt{\cosh \sqrt{z}}, \quad g(z) = \sqrt{(z-i)(z^2-1)}, \quad h(z) = \sqrt{z^4-1}$$

of the complex variable z .

- Choose branch cuts for f to be single valued in the complex z plane.
- Can the function g be defined to be single valued in the domain exterior to the circle $|z| = 2$? Discuss.
- Answer part (b) for the function h , and explain.

- (5) Find the first term in the asymptotic expansion of

$$f(x) = \int_C \frac{1}{\zeta^3} \exp \left[i x \left(\zeta^2 - \frac{1}{\zeta^2} \right) \right] d\zeta$$

as $x \rightarrow +\infty$, where C is a closed contour around the origin.

- (6) Consider the rapidly varying diffusivity

$$K(x, \epsilon) = A + G(x/\epsilon)$$

where A is chosen to guarantee K is positive, and ϵ is a small constant. Also consider the function, $v(x)$. By applying homogenization, average the following advection-diffusion equation

$$\frac{\partial u}{\partial t} + \frac{1}{\epsilon} v(x/\epsilon) \frac{\partial u}{\partial y} = \frac{\partial}{\partial x} \left(K(x, \epsilon) \frac{\partial u}{\partial x} \right) + \frac{\partial}{\partial y} \left(K(x, \epsilon) \frac{\partial u}{\partial y} \right)$$

$$u(x, y, 0) = u_0(x, y),$$

and calculate a leading order effective equation governing the evolution as $\epsilon \rightarrow 0$ over the (x, y) -plane, assuming the functions $G(x/\epsilon)$, $v(x/\epsilon)$ are mean zero, periodic, and share the same period. Solve the averaged equation in free space. If instead we had $v(x/\epsilon^2)$, how would things change?

- (7) (a) Find two term asymptotic expansions as $\epsilon \rightarrow 0$ for roots of the equation

$$\epsilon z^7 + (z - 5)^2 = \epsilon z$$

(analyze carefully the unperturbed root at $z = 5$, and analyze one root coming in from infinity.)

- (b) Find a two term asymptotic expansion for the $x = 0$ unperturbed root of the equation

$$x^2 e^{-x^4} = \epsilon$$

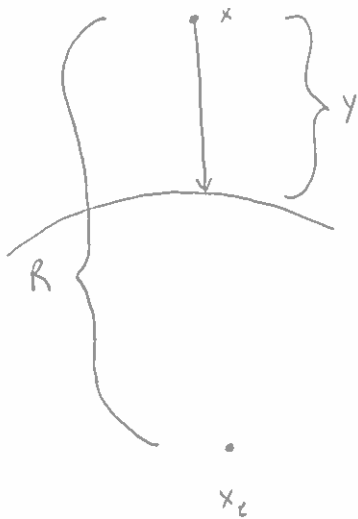
- (8) (a) Explain the difference between pointwise convergence and asymptotic convergence. Illustrate with the particular example of power series.

- (b) Given a real integral, $\int_a^b f(x, y) dy$, with $f(x, y) \sim f_0(y)$ as $x \rightarrow x_0$, under what conditions is termwise integration guaranteed to yield valid asymptotics?

Problem 1

Aug 2017

a)



$$m \ddot{x} = - \frac{GmM}{|x-x_c|^2} - \beta \dot{x}^3$$

$$M \ddot{x}_c = \frac{GmM}{|x-x_c|^2}$$

since $M \gg m$, $\dot{x}_c \ll \dot{x}$ b/c the velocity of an object is inversely proportional to its mass. Thus $\dot{x}_c \approx \text{const}$ so $\ddot{x}_c \approx 0$.

$$\begin{cases} \ddot{x} = - \frac{GM}{|x-x_c|^2} - \frac{\beta}{m} \dot{x}^3 \\ \dot{x}(0) = 0 \\ x(0) = y_0 + R_c \quad y_0 > 0 \end{cases}$$

change of var:

$$\begin{cases} y = x - R_c \Rightarrow x = y + R_c \\ x_c = 0 \end{cases} \quad \begin{matrix} \uparrow \\ \text{radius of earth} \end{matrix}$$

$$\Rightarrow \ddot{y} = - \frac{GM}{|y+R_c|^2} - \frac{\beta}{m} \dot{y}^3$$

$$\boxed{\ddot{y} = - \underbrace{\frac{GM}{R_c^2}}_g \frac{1}{\left(1 + \frac{y}{R_c}\right)^2} - \frac{\beta}{m} \dot{y}^3}$$

where $y(0) = x(0) - R_c = y_0 > 0$

$$\dot{y}(0) = \dot{x}(0) = 0$$

b) Let $v_t =$ terminal velocity and $L = \frac{v_t^2}{g}$, $T = \frac{v_t}{g}$.

Choose $y = Lz$ and $t = T\tau \rightarrow \frac{dt}{d\tau} = T$
non-dim

$$\frac{dy}{dt} = \frac{d(Lz)}{d\tau} \frac{d\tau}{dt} = \frac{L}{T} \frac{dz}{d\tau} = v_t \dot{z}$$

$$\frac{d^2y}{dt^2} = \frac{d}{dt} \left(\frac{L}{T} \frac{dz}{d\tau} \right) = \frac{L}{T^2} \frac{d^2z}{d\tau^2} = g \ddot{z}$$

$$\Rightarrow g \ddot{z} = -g \frac{1}{\left|1 + \frac{zL}{R_c}\right|^2} - \frac{\rho}{m} (v_t \dot{z})^3$$

$$\text{Let } \varepsilon = \frac{L}{R_c} \text{ and } \alpha = \frac{\rho v_t R_c}{m}$$

$$\ddot{z} = - \frac{1}{\left|1 + \varepsilon z\right|^2} - \alpha \dot{z}^3 \quad (1)$$

$$0 = \dot{y}(0) = v_t \dot{z}(0)$$

$$\dot{z}(0) = 0$$

$$y_0 = y(0) = Lz(0)$$

$$z(0) = \frac{y_0}{L} = \bar{z}$$

c) We know $v_t^2 \ll R \varepsilon \cdot g$ so $\varepsilon \ll 1$. Now consider

$$z \sim z_0 + \varepsilon z_1 + \dots$$

and plug into (1):

$$\ddot{z}_0 + \varepsilon \ddot{z}_1 = -\frac{1}{(1 + \varepsilon(z_0 + \varepsilon z_1 + \dots))} - \varepsilon \alpha (z_0 + \varepsilon z_1 + \dots)^3$$

$$= -1 + 2\varepsilon(z_0 + \varepsilon z_1 + \dots) + \mathcal{O}(\varepsilon^2) - \varepsilon \alpha (z_0 + \varepsilon z_1 + \dots)^3$$

$$\mathcal{O}(1): \ddot{z}_0 = -1$$

$$\dot{z}_0 = -\tau + \alpha \rightarrow 0$$

$$z_0 = -\frac{\tau^2}{2} + b = -\frac{1}{2}\tau^2 + \bar{z}$$

$$\mathcal{O}(\varepsilon): \ddot{z}_1 = 2z_0 - \alpha \dot{z}_0 = 2(-\frac{1}{2}\tau^2 + \bar{z}) + \alpha \tau = -\tau^2 + 2\alpha\tau + \bar{z}$$

$$\dot{z}_1 = -\frac{1}{3}\tau^3 + \frac{\alpha}{2}\tau^2 + \bar{z}\tau + \alpha \rightarrow 0$$

$$z_1 = -\frac{1}{12}\tau^4 + \frac{\alpha}{6}\tau^3 + \frac{\bar{z}}{2}\tau^2 + b \rightarrow 0$$

$$z \sim -\frac{1}{2}\tau^2 + \bar{z} + \varepsilon \left(-\frac{1}{12}\tau^4 + \frac{\alpha}{6}\tau^3 + \frac{\bar{z}}{2}\tau^2 \right)$$



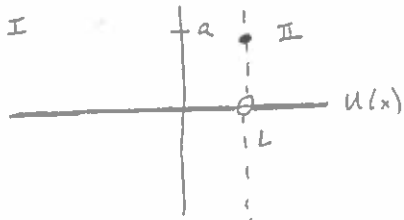
Problem 2

Aug 2017

$$\epsilon y'' - (u(x) + \lambda) y = 0 \quad \Rightarrow \quad y'' = \frac{\lambda + u(x)}{\epsilon} y$$

$$u(x) = a \delta(x-L) + b \delta(x+L)$$

1) $b=0$



Since we want $y \rightarrow 0$ as $|x| \rightarrow \infty$, we know $\frac{\lambda + u(x)}{\epsilon} > 0$ in regions I and II. So $\frac{\lambda}{\epsilon} > 0$.

Region I & II $y_i'' = \underbrace{\frac{\lambda}{\epsilon}}_{> 0} y_i \quad i=1,2$

$$y_i = A_i e^{\sqrt{\frac{\lambda}{\epsilon}} x} + B_i e^{-\sqrt{\frac{\lambda}{\epsilon}} x}$$

Since $y \rightarrow 0$ as $|x| \rightarrow \infty$, we know $B_1 = 0$ and $A_2 = 0$

$$\Rightarrow y_1 = A e^{\sqrt{\frac{\lambda}{\epsilon}} x}$$

$$y_2 = B e^{-\sqrt{\frac{\lambda}{\epsilon}} x}$$

Patching:

$$y_1(L) = y_2(L)$$

$$A e^{-L\sqrt{\lambda/\epsilon}} = B e^{-L\sqrt{\lambda/\epsilon}}$$

$$B = A e^{2L\sqrt{\lambda/\epsilon}}$$

Let $0 < \gamma \ll 1$

$$\int_{L-\gamma}^{L+\gamma} y'' = \int_{L-\gamma}^{L+\gamma} \frac{\lambda + u(x)}{\epsilon} y = \int_{L-\gamma}^{L+\gamma} \frac{\lambda}{\epsilon} y + \int_{L-\gamma}^{L+\gamma} \frac{a}{\epsilon} y \delta(x-L)$$

$\rightarrow 0$ as $\gamma \rightarrow 0$

$$y' \Big|_{L-\gamma}^{L+\gamma} = \frac{a}{\epsilon} y(L)$$

as $\gamma \rightarrow 0$: $y_2'(L) - y_1'(L) = \frac{a}{\epsilon} y(L)$

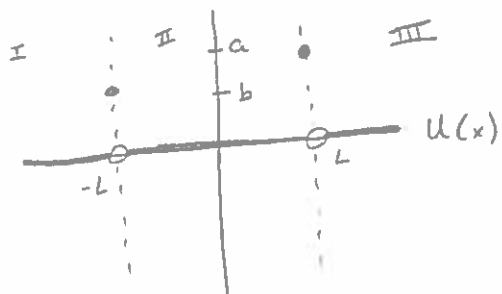
$$-B\sqrt{\lambda/\epsilon} e^{-L\sqrt{\lambda/\epsilon}} - A\sqrt{\lambda/\epsilon} e^{-L\sqrt{\lambda/\epsilon}} = \frac{a}{\epsilon} A e^{-L\sqrt{\lambda/\epsilon}}$$

$$-A\sqrt{\lambda/\epsilon} e^{\sqrt{\lambda/\epsilon}} - A\sqrt{\lambda/\epsilon} e^{-L\sqrt{\lambda/\epsilon}} = \frac{a}{\epsilon} A e^{-L\sqrt{\lambda/\epsilon}}$$

$$2\sqrt{\lambda/\epsilon} = -\frac{a}{\epsilon}$$

$$\frac{\lambda}{\epsilon} = \frac{a^2}{4\epsilon^2}$$

$$\boxed{\lambda = \frac{a^2}{4\epsilon}}$$

2) $b \neq 0$ 

Again, $y \rightarrow 0$ as $|x| \rightarrow \infty$ so $\frac{\lambda}{\varepsilon} > 0$.

All Regions: $i=1, 2, 3$

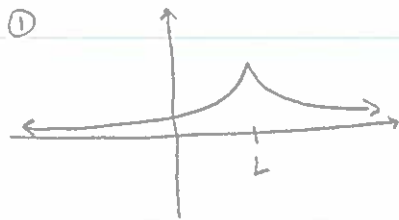
$$y_i = A_i e^{\sqrt{\frac{\lambda}{\varepsilon}} x} + B_i e^{-\sqrt{\frac{\lambda}{\varepsilon}} x}$$

where $A_3 = 0$ and $B_1 = 0$ to ensure $y \rightarrow 0$ as $|x| \rightarrow \infty$

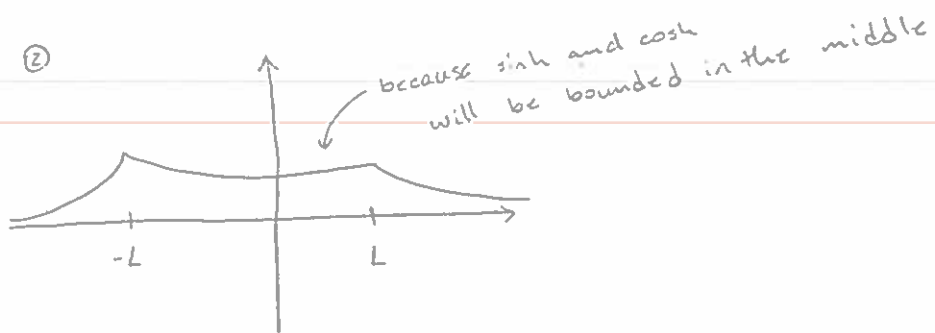
$$\Rightarrow \begin{cases} y_1 = A_1 e^{\sqrt{\frac{\lambda}{\varepsilon}} x} \\ y_2 = A_2 e^{\sqrt{\frac{\lambda}{\varepsilon}} x} + B_2 e^{-\sqrt{\frac{\lambda}{\varepsilon}} x} \\ y_3 = B_3 e^{-\sqrt{\frac{\lambda}{\varepsilon}} x} \end{cases}$$

To find eigenvalues, one must proceed with a patching process similar to the one in part 1.

The result from part 1 looks like:



While the result in part 2 looks like:

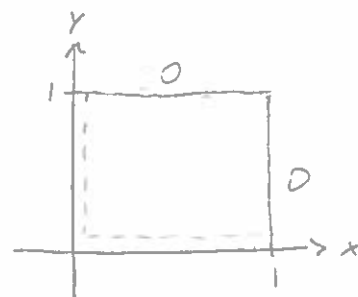


If $L \rightarrow \infty$, you can see solu ① will stretch out the left portion and essentially y will be zero. But solu ② will have the middle section stretched out and it will more closely represent the constant line $y = \frac{a+b}{2}$.

Problem 3

Aug 2017

$$\begin{cases} \varepsilon u_{xx} + u_{yy} = \varepsilon u^2 \\ u(0, y) = g(y) \quad , \quad u(x, 0) = f(x) \quad \text{outer BC's} \\ u(1, y) = 0 \quad , \quad u(x, 1) = 0 \quad \text{inner} \end{cases}$$



→ $\mathcal{O}(1)$

Outer: $u_{yy} = 0$

$$u(x, y) = a(x)y + b(x)$$

$$u(0, y) = a(0)y + b(0) = g(y)$$

$$u(x, 0) = b(x) = f(x)$$

Inner:

$$x=1: \quad z = \frac{x-1}{\varepsilon} \Rightarrow \frac{d}{dx} = \frac{1}{\varepsilon} \frac{d}{dz}$$

$$\frac{1}{\varepsilon} v_{xx} + v_{yy} = \varepsilon v^2$$

$$y=1: \quad z = \frac{y-1}{\varepsilon} \Rightarrow \frac{d}{dy} = \frac{1}{\varepsilon} \frac{d}{dz}$$



Problem 4

a) $f(z) = \sqrt{\cosh(\sqrt{z})}$

Branch points:

1) $z=0$

2) $\cosh(\sqrt{z}) = 0$

$$\cos(i\sqrt{z}) = 0$$

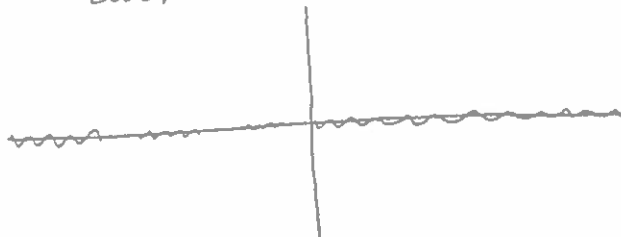
$$\cos(\sqrt{-z}) = 0$$

$$\sqrt{-z} = \frac{\pi}{2} + \pi n$$

$$-z = \left(\frac{\pi}{2} + \pi n\right)^2$$

$$z = -\left(\frac{\pi}{2} + \pi n\right)^2 = -\frac{\pi^2}{4}, -\frac{9\pi^2}{4}, -\frac{25\pi^2}{4}, \dots$$

Branch cut:



b) $g(z) = \sqrt{(z-i)(z^2-1)}$

Branch points: $z = i, \pm 1$

One of the branch cuts will have to go to infinity to ensure $g(z)$ is single valued.

By excluding $|z| \leq 2$, you lose the branch points you'd need to make the branch cut so $g(z)$ can't be s.v. on this domain.

c) $h(z) = \sqrt{z^4 - 1}$

Branch points: $z = \pm i, \pm 1$

All branch cuts are made on the interior of $|z| \leq 2$ to make $h(z)$ s.v. so removing $|z| \leq 2$ from the $h(z)$ will not cause any problems.

Problem 5

Aug 2017

$$f(x) = \int_C \frac{1}{s^3} e^{ix(s^2 - \frac{1}{s^2})} ds$$

No branch cuts, just a 3rd order pole at $s=0$.

$$\text{Let } s = e^{i\theta} \Rightarrow ds = -ie^{i\theta} d\theta \text{ where } \theta \in [0, 2\pi)$$

$$f(x) = i \int_0^{2\pi} \frac{1}{e^{3i\theta}} e^{ix(e^{2\theta i} - e^{-2\theta i})} e^{i\theta} d\theta \quad \sin \theta = \frac{e^{i\theta} - e^{-i\theta}}{2i}$$

$$= i \int_0^{2\pi} e^{-2\theta i} e^{-2x \sin(2\theta)} d\theta$$

$$= i \int_0^{2\pi} \cos(2\theta) e^{-2x \sin(2\theta)} d\theta + \int_0^{2\pi} \sin(2\theta) e^{-2x \sin(2\theta)} d\theta$$

Now the function is defined in \mathbb{R} and we are integrating in \mathbb{R} so we can use Laplace's Method.

$$g_1(\theta) = \cos(2\theta)$$

$$g_2(\theta) = \sin(2\theta)$$

$$h(\theta) = -2 \sin(2\theta)$$

$$h'(\theta) = -4 \cos(2\theta) = 0$$

$$\theta^* = \frac{\pi}{4}, \frac{3\pi}{4}, \frac{5\pi}{4}, \frac{7\pi}{4}$$

$$h\left(\frac{\pi}{4}\right) = h\left(\frac{3\pi}{4}\right) = -2 \sin\left(\frac{\pi}{2}\right) = -2$$

$$\text{dominant} \rightarrow h\left(\frac{3\pi}{4}\right)_{\theta_1} = h\left(\frac{7\pi}{4}\right)_{\theta_2} = -2 \sin\left(\frac{3\pi}{2}\right) = 2$$

θ_1 and θ_2 have equal weights

$$h''(\theta) = 8 \sin(2\theta)$$

$$h''(\theta_1) = 8 \sin\left(\frac{3\pi}{2}\right) = -8$$

$$f(x) \sim \sum_{j=1}^2 (ig_1(\theta_j) + g_2(\theta_j)) e^{xh(\theta_j)} \sqrt{\frac{2\pi}{x|h''(\theta_j)|}} \quad \text{as } x \rightarrow \infty$$

$$g_1(\theta_1) = \cos\left(\frac{3\pi}{2}\right) = 0$$

$$g_2(\theta_1) = \sin\left(\frac{3\pi}{2}\right) = -1$$

$$g_1(\theta_2) = 0$$

$$g_2(\theta_2) = -1$$

$$f(x) \sim -2 e^{2x} \sqrt{\frac{\pi}{4x}} \quad \text{as } x \rightarrow \infty$$

$$f(x) = -e^{2x} \sqrt{\frac{\pi}{x}} \quad \text{as } x \rightarrow \infty$$

$$\begin{cases} u_t + \frac{1}{\varepsilon} v\left(\frac{x}{\varepsilon}\right) \cdot \partial_x (k \partial_x u) + \partial_y (k \partial_y u) \\ u(x, y, 0) = u_0(x, y) \end{cases}$$

where $k(x, \varepsilon) = A + G\left(\frac{x}{\varepsilon}\right)$.

Let $z = \frac{x}{\varepsilon}$, then $\partial_x \mapsto \partial_x + \frac{1}{\varepsilon} \partial_z$

Ansatz ε : $u = \bar{u} + \varepsilon u_1 + \dots$

$$u_t = (\partial_x + \frac{1}{\varepsilon} \partial_z) (k (\partial_x + \frac{1}{\varepsilon} \partial_z) u) + \partial_y (k \partial_y u) - \frac{1}{\varepsilon} v(z) \partial_y u$$

$O(\frac{1}{\varepsilon^2})$: $\partial_z (k \partial_z \bar{u}) = 0 \Rightarrow \bar{u}$ indep z

$O(\frac{1}{\varepsilon})$: $\partial_z (k \partial_z u_1) + \partial_z (k \partial_x \bar{u}) + \frac{\partial_x (k \partial_z \bar{u})}{\bar{u} \text{ indep } z} - v(z) \partial_y \bar{u} = 0$

$\langle \partial_z [k (\partial_z u_1 + \partial_x \bar{u})] \rangle_z - \langle \frac{v(z) \partial_y \bar{u}}{\substack{\uparrow \\ \text{periodic}}} \rangle_z = 0$
 \uparrow indep z

$k (\partial_z u_1 + \partial_x \bar{u}) = A$ indep z

~~$\langle \partial_z u_1 \rangle_z$~~ , $\langle \partial_x \bar{u} \rangle_z = \left\langle \frac{A}{k} \right\rangle_z$
 FA

$\partial_x \bar{u} = \frac{A}{\langle k \rangle_z} \Rightarrow A = \langle k \rangle_z \partial_x \bar{u}$

$k (\partial_z u_1 + \partial_x \bar{u}) = \langle k \rangle_z \partial_x \bar{u}$

Also, u_1 indep of z \therefore RHS is

$$\Theta(1): \partial_z (k \partial_z u_2) + \partial_z (k \partial_x u_1) + \partial_x (k \partial_z u_1) - v(z) \partial_y u_1$$

$$+ \partial_x (k \partial_x \bar{u}) + \partial_y (k \partial_y \bar{u}) = \bar{u}_t$$

$$\underbrace{\langle \partial_z (k (\partial_z u_2 + \partial_x u_1)) \rangle_z}_0 - \underbrace{\langle v(z) \partial_y u_1 \rangle_z}_0$$

\downarrow periodic \downarrow indep z

FA

$$+ \langle \partial_x (k (\partial_z u_1 + \partial_x \bar{u})) \rangle_z + \langle \partial_y (k \partial_y \bar{u}) \rangle_z = \bar{u}_t$$

$\underbrace{\langle k \rangle_z}_\alpha \partial_x \bar{u}$

$$\bar{u}_t = \underbrace{\langle k \rangle_z}_\alpha \bar{u}_{xx} + \underbrace{\langle k \rangle_z}_\beta \bar{u}_{yy}$$

$$\bar{u} = \mathcal{F}^{-1} \left\{ e^{(\alpha \chi^2 + \beta \eta^2) t} \right\}$$

where $\mathcal{F}(\bar{u}_{xx}) = \chi^2$

and $\mathcal{F}(\bar{u}_{yy}) = \eta^2$

a) $\varepsilon z^7 + (z-5)^2 = \varepsilon z$ (1)

Unperturbed root:

$$(z-5)^2 = 0$$

$$z = 5$$

$$\Rightarrow z_0 = \sum_{n=0}^{\infty} a_n \varepsilon^{n/2} \quad \text{where } a_0 = 5$$

Plug z_0 into (1):

$$\varepsilon (5 + a_1 \varepsilon^{1/2})^7 + (a_1 \varepsilon^{1/2})^2 = \varepsilon (5 + a_1 \varepsilon^{1/2})$$

$$O(\varepsilon): 5^7 + a_1^2 = 5$$

$$a_1 = \pm \sqrt{5 - 5^7} = \pm i \sqrt{5^7 - 5}$$

$$\Rightarrow \begin{aligned} z_1 &\sim 5 + i \sqrt{5^7 - 5} \varepsilon^{1/2} \\ z_2 &\sim 5 - i \sqrt{5^7 - 5} \varepsilon^{1/2} \end{aligned}$$

Let $x = \varepsilon^\alpha z \Rightarrow z = \varepsilon^{-\alpha} x$ and plug into (1)

$$\varepsilon^{1-7\alpha} x^7 + (\varepsilon^{-\alpha} x - 5)^2 = \varepsilon^{1-\alpha} x$$

$$\varepsilon^{1-7\alpha} x^7 + \varepsilon^{-2\alpha} x^2 - 10 \varepsilon^{-\alpha} x + 25 = \varepsilon^{1-\alpha} x$$

$$1 - 7\alpha = -2\alpha$$

$$\alpha = 1/5$$

$$x^7 + x^2 - 10 \varepsilon^{1/5} x + 25 \varepsilon^{2/5} - \varepsilon^{4/5} x = 0$$

(2)

Consider the unperturbed version of (2):

$$x^7 + x^2 = 0$$

$$x^2(x^5 + 1) = 0$$

$$\omega = -\epsilon^{i \frac{2\pi}{5}}$$

$$x = 0, 0, \omega, \omega^2, \omega^3, \omega^4, -1$$

↑ correspond to z_1 and z_2

$$x_{k+2} = \omega^k + \sum_{n=1}^{\infty} b_n^k \epsilon^{n/5} \quad k = 1, \dots, 5$$

Plug x_{k+2} into (2)

$$\begin{aligned} (\omega^k + b_1^k \epsilon^{1/5})^7 + (\omega^k + b_1^k \epsilon^{1/5})^2 - 10 \epsilon^{1/5} (\omega^k + b_1^k \epsilon^{1/5}) + 25 \epsilon^{2/5} \\ - \epsilon^{4/5} (\omega^k + b_1^k \epsilon^{1/5}) = 0 \end{aligned}$$

$$\theta(1): \omega^{7k} + \omega^{2k} = 0 \quad \checkmark$$

$$\theta(\epsilon^{1/5}): \omega^{6k} b_1^k + 2\omega^k b_1^k - 10\omega^k = 0$$

$$b_1^k = \frac{10}{\omega^{5k} + 2}$$

$$b_1^k = \frac{10}{-\omega^k + 2}$$

Recall: $z_j = x_j \epsilon^{-j/5}$

$$z_3 \sim \frac{\omega}{\epsilon^{1/5}} + \frac{10}{2 - \omega}$$

$$z_4 \sim \frac{\omega^2}{\epsilon^{1/5}} + \frac{10}{2 - \omega^2}$$

⋮

$$z_1 \sim -\frac{1}{\epsilon^{1/5}} + \frac{10}{3}$$

$$b) \quad x^2 = \varepsilon e^{x^4} \quad (1)$$

$$x^2 = \varepsilon(1 + x^4) \quad (2)$$

$$\text{Unperturbed: } x^2 = 0$$

$$x_0 = 0$$

$$\Rightarrow x_0 = \sum_{n=1}^{\infty} a_n \varepsilon^{n/2} = a_1 \varepsilon^{1/2} + a_2 \varepsilon + a_3 \varepsilon^{3/2} + \dots$$

Plug x_0 into (2) and find a hierarchy of eqns:

$$\mathcal{O}(1): 0 = 0$$

$$\mathcal{O}(\varepsilon^{1/2}): 0 = 0$$

$$\mathcal{O}(\varepsilon): a_1^2 = 1$$

$$a_1 = \pm 1$$

$$\mathcal{O}(\varepsilon^{3/2}): 2a_1 a_2 = 0$$

$$a_2 = 0$$

$$\mathcal{O}(\varepsilon^2): 2a_1 a_3 + a_2^2 = 0$$

$$a_3 = 0$$

$$\mathcal{O}(\varepsilon^{5/2}): 2a_1 a_4 + \cancel{2a_2 a_3}^0 = 0$$

$$a_4 = 0$$

$$\mathcal{O}(\varepsilon^3): \cancel{2a_2 a_4}^0 + 2a_1 a_5 + \cancel{a_3^2}^0 = a_1^4$$

$$a_5 = \frac{a_1^3}{2}$$

$$x_1 \sim \varepsilon^{1/2} + \frac{1}{2} \varepsilon^{5/2}$$

$$x_2 \sim -\varepsilon^{1/2} - \frac{1}{2} \varepsilon^{5/2}$$



a) Let $\Omega \subseteq \mathbb{C}$, $x_0 \in \Omega$ and f_n be a sequence of fncs.

Ptws: We say $f_n \rightarrow f$ conv ptws if $\forall \varepsilon > 0 \exists N \in \mathbb{N}$ st
 $n > N \Rightarrow |f_n(x_0) - f(x_0)| < \varepsilon$.

Note: we are only concerned with fixed x_0 and N can vary.

Asymp: We say $f_n \rightarrow f$ conv asymp'ly if $\forall \varepsilon > 0$
 $\exists \delta > 0$ st $|x - x_0| < \delta \Rightarrow |f_n(x) - f(x)| < \varepsilon$.
 for fixed N

Note: This time N is fixed and x can vary.

Ex: In each example I'll use ratio test to show conv.

i) $\sum_{n=0}^{\infty} \frac{x^n}{n!}$

$$\left| \frac{x^{n+1}}{(n+1)!} \cdot \frac{n!}{x^n} \right| = \left| \frac{x}{n+1} \right| \xrightarrow[n \rightarrow \infty]{x \text{ fixed}} 0 \quad \therefore \text{ptws conv}$$

$$\xrightarrow[n \text{ fixed}]{x \rightarrow \infty} \infty \quad \therefore \text{not asymp conv}$$

ii) $\sum_{n=0}^{\infty} \frac{n!}{x^n}$

$$\left| \frac{(n+1)!}{x^{n+1}} \cdot \frac{x^n}{n!} \right| = \left| \frac{n+1}{x} \right| \xrightarrow[n \rightarrow \infty]{x \text{ fixed}} \infty \quad \therefore \text{not ptws conv}$$

$$\xrightarrow[n \text{ fixed}]{x \rightarrow \infty} 0 \quad \therefore \text{asymp conv}$$

b) When f is uniformly continuous and integrable near x_0 .